

Storage of a single photon in a Bose-Einstein condensate

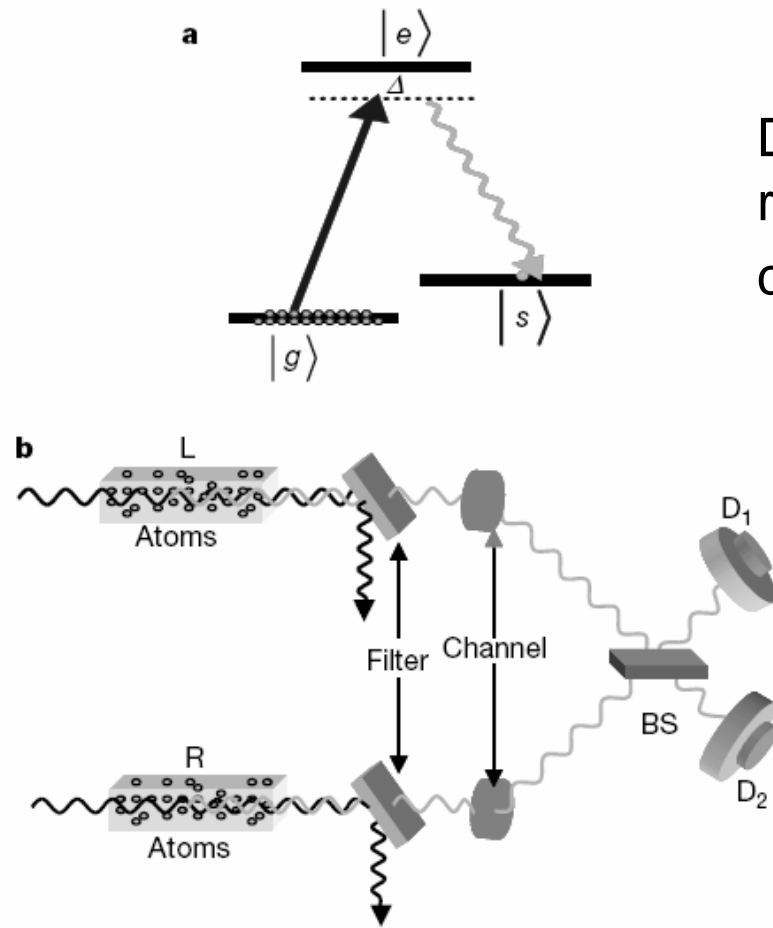
December 13, 2006

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Motivation: DLCZ protocol



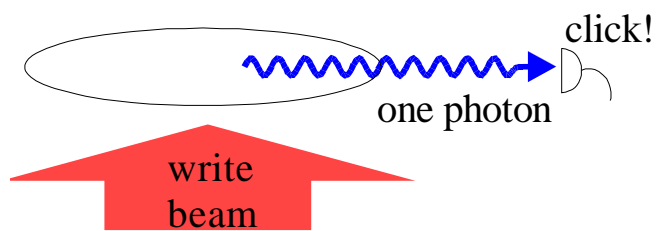
Detection of a forward-scattered photon results in the excitation of the symmetric collective mode defined by

$$S^+ \equiv \frac{1}{\sqrt{N}} \sum_{i=1}^N |s\rangle_i \langle g|$$

L.-M. Duan, M. D. Lukin, J. I. Cirac, and P. Zoller, Nature. **414**, 413 (2001)

Writing, storing, and reading of a single photon

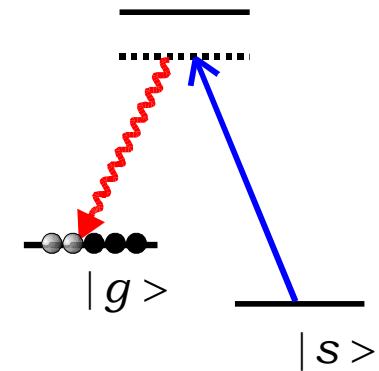
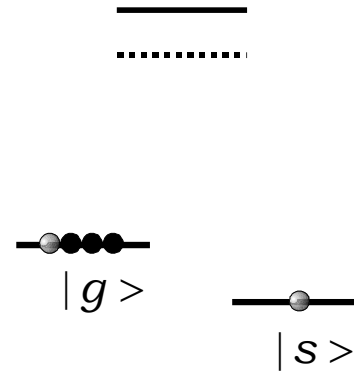
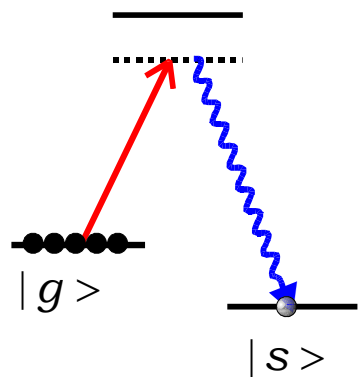
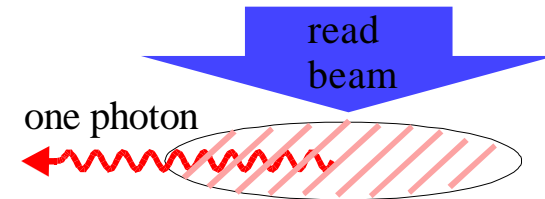
writing



storing



reading

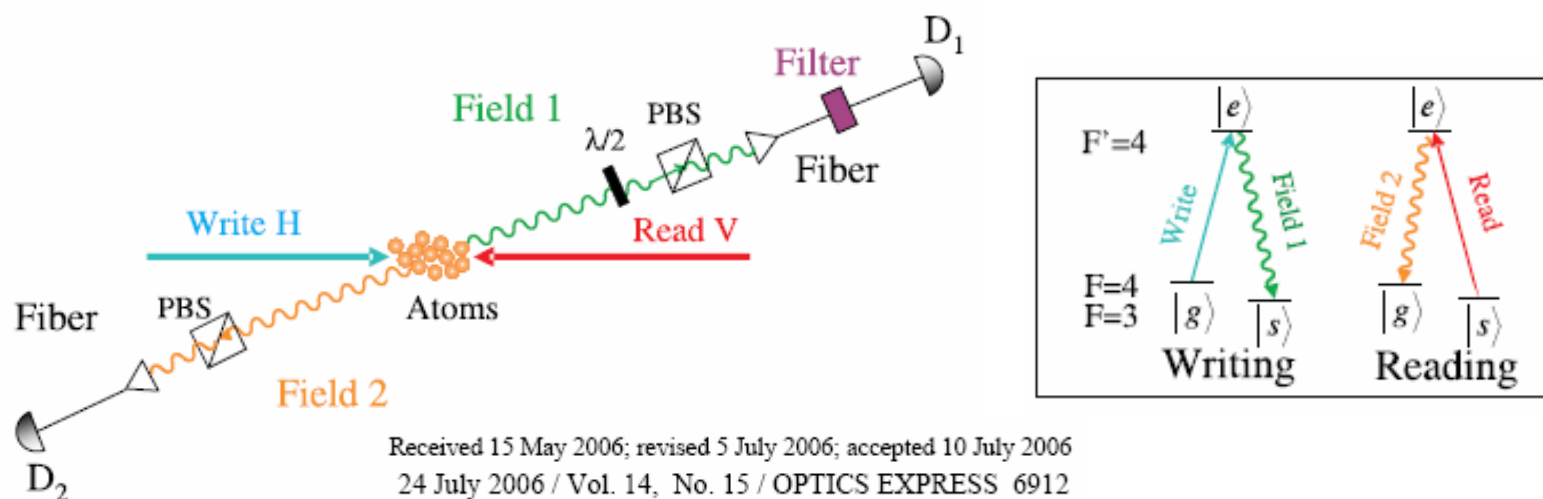


Efficient retrieval of a single excitation stored in an atomic ensemble

Julien Laurat, Hugues de Riedmatten, Daniel Felinto, Chin-Wen Chou, Erik W. Schomburg, and H. Jeff Kimble

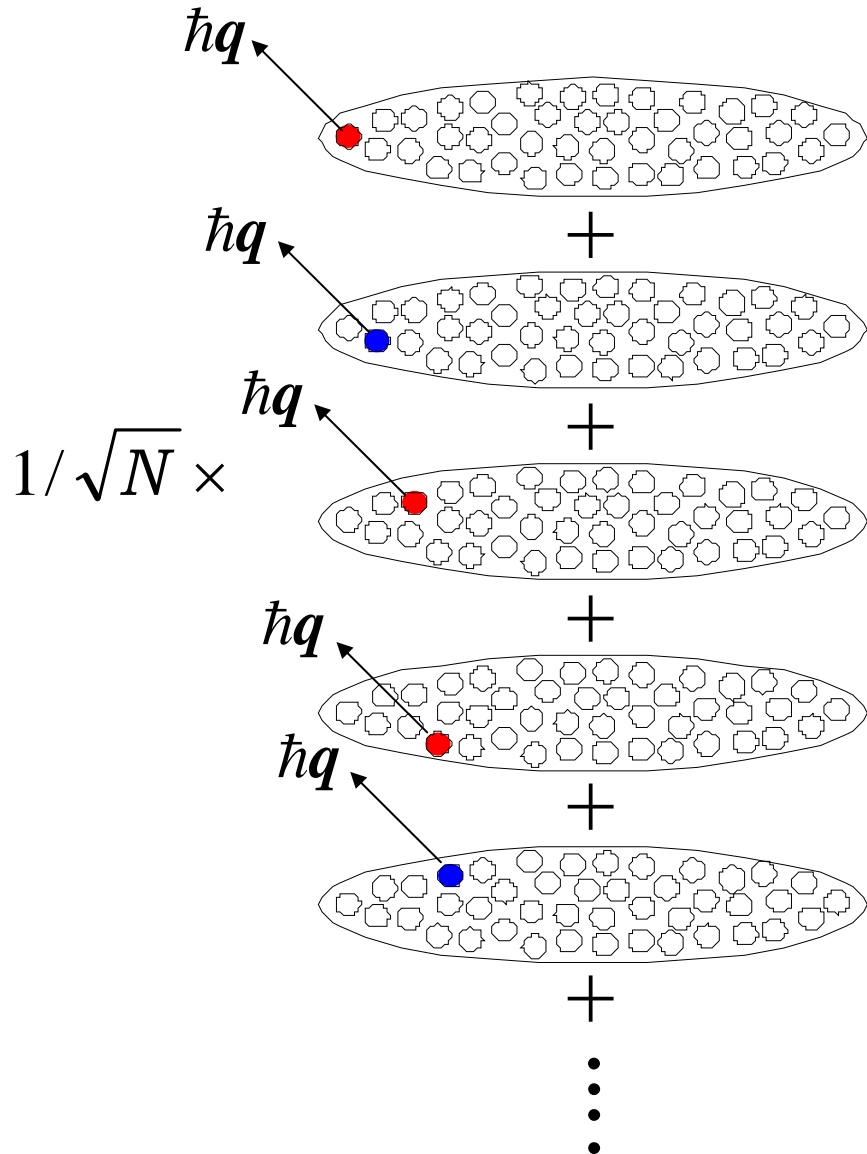
Norman Bridge Laboratory of Physics 12-33, California Institute of Technology, Pasadena, California 91125, USA

Abstract: We report significant improvements in the retrieval efficiency of a single excitation stored in an atomic ensemble and in the subsequent generation of strongly correlated pairs of photons. A 50% probability of transforming the stored excitation into one photon in a well-defined spatio-temporal mode at the output of the ensemble is demonstrated. These improvements are illustrated by the generation of high-quality heralded single photons with a suppression of the two-photon component below 1% of the value for a coherent state. A broad characterization of our system is performed for different parameters in order to provide input for the future design of realistic quantum networks.

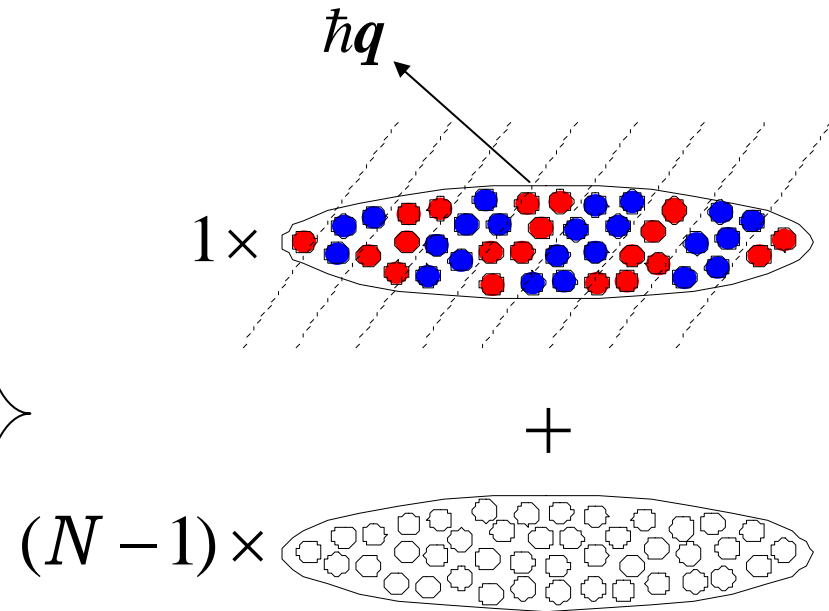


Received 15 May 2006; revised 5 July 2006; accepted 10 July 2006
24 July 2006 / Vol. 14, No. 15 / OPTICS EXPRESS 6912

The origin of a grating (Indiscernability of the atoms)



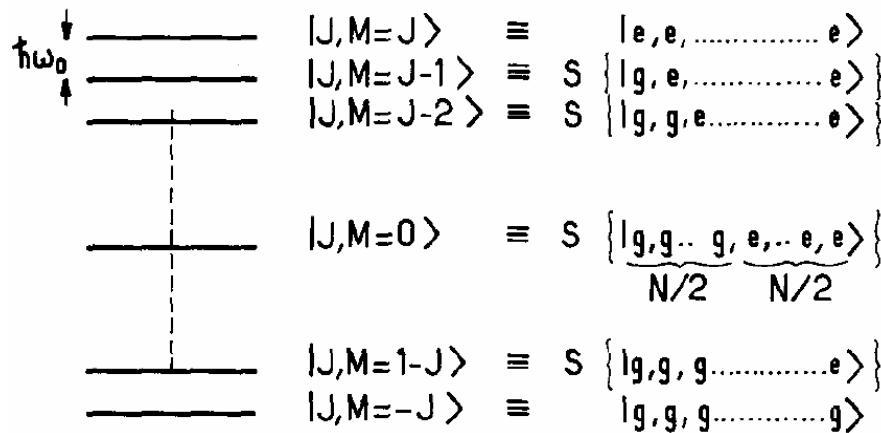
One atom is excited to the collective atomic mode defined by S^+



$$S^+ |0_a\rangle \equiv \frac{1}{\sqrt{N}} \sum_{i=1}^N |g_1, g_2, \dots, s_i, \dots, g_{N_a}\rangle$$

Collective mode = Dicke state

N-atom system \Leftrightarrow N spin-1/2 system with the total spin $J = N/2$
 (assumption: *Indiscernability* of the atoms with respect to photon emission)



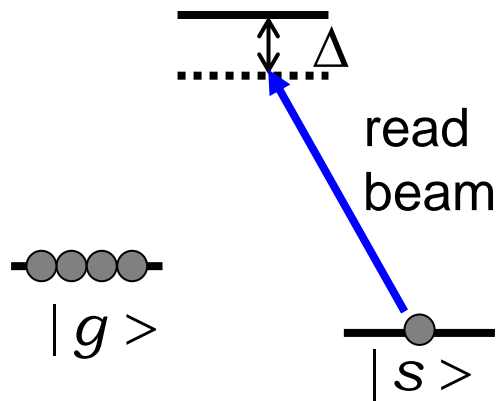
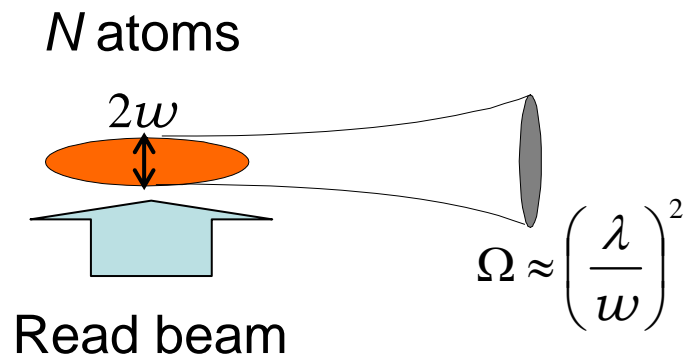
Spontaneous emission rate of the N-atom system:

$$\begin{aligned} \Gamma_N &= \Gamma \langle J, M | J_+ J_- | J, M \rangle \\ &= \Gamma (J + M)(J - M + 1) \\ &= \Gamma N_e (N_g + 1) \end{aligned}$$

Bosonic stimulation
 (Superradiant emission)

R. H. Dicke, Phys. Rev. **93**, 99 (1954)
 M. Gross and S. Haroche, Phys. Rep. **93**, 301 (1982)

Raman scattering rate for a cigar-shaped atomic ensemble



Single-atom Raman scattering rate

$$R = \Gamma \frac{\Omega_p^2}{4\Delta^2}$$

N_a -atom Raman scattering rate

$$R_N = f(\theta) R \langle J, M | J_+ J_- | J, M \rangle \Omega$$

$$= f(\theta) R N_s (N_g + 1) \Omega$$

Mode field pattern

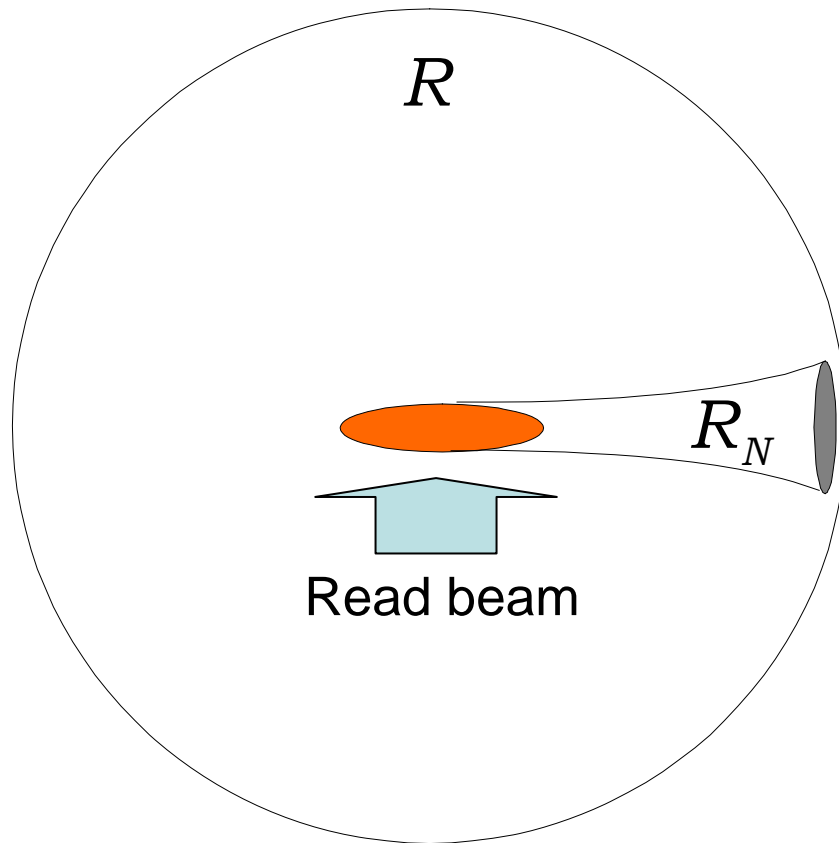
Phase matching solid angle

For the reading ($N_g = N-1$, $N_s = 1$)

$$R_N = \boxed{N} f(\theta) R \Omega$$

Collective enhancement

Spontaneous scattering vs. Collective scattering



The ratio between spontaneous and collective Raman scattering rates:

$$\frac{R_N}{R} = N\eta \quad : \text{ cooperativity parameter}$$

$$\eta \equiv f(\theta)\Omega \quad : \text{ single-atom optical depth}$$

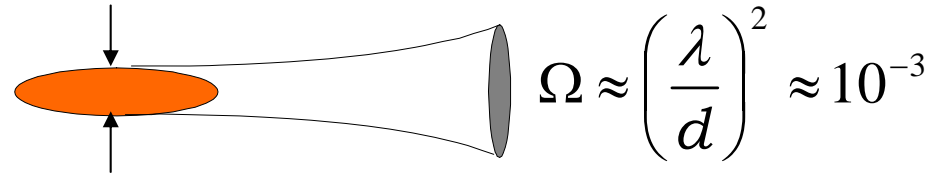
$$\Omega \approx \left(\frac{\lambda}{w}\right)^2 f(\theta) = \begin{cases} \frac{3 \sin^2 \theta}{8\pi} & (\pi\text{-pol.}) \\ \frac{3(1 + \cos^2 \theta)}{16\pi} & (\sigma\text{-pol.}) \end{cases}$$

The probability that an atom in the collective mode emits a photon into the solid angle Ω

$$P_s = \frac{N\eta}{1 + N\eta}$$

Cooperativity parameter of Bose condensates

Typical size of a Bose condensate: $d = 10 \mu\text{m}$



Typical number of atoms in a Bose condensate: $N = 10^6$

Cooperativity parameter for a typical Bose condensate:

$$N\eta \approx N\Omega \approx 10^3$$

Probability of successful retrieval of a single photon:

$$P_s = \frac{N\eta}{1 + N\eta} \approx 99.9\%$$

BEC is ideal for storage of a single photon!

cf.) Cooperativity parameter and Purcell factor for cavities

The rate for an excited atom in the cavity to emit a photon into the cavity mode

$$R = \frac{2\pi}{\hbar^2} |\langle g, 1 | \hbar g_0 (a\sigma^+ + a^\dagger\sigma) | e, 0 \rangle|^2 \delta(\omega - \omega_A)$$

$$= 2\pi g_0^2 \frac{k/\pi}{\kappa^2 + \delta^2} \xrightarrow{\delta=0} \frac{2g_0^2}{\kappa} \quad \frac{k/\pi}{\kappa^2 + \delta^2} \begin{array}{l} \text{Normalized} \\ \text{Cavity line shape} \end{array}$$

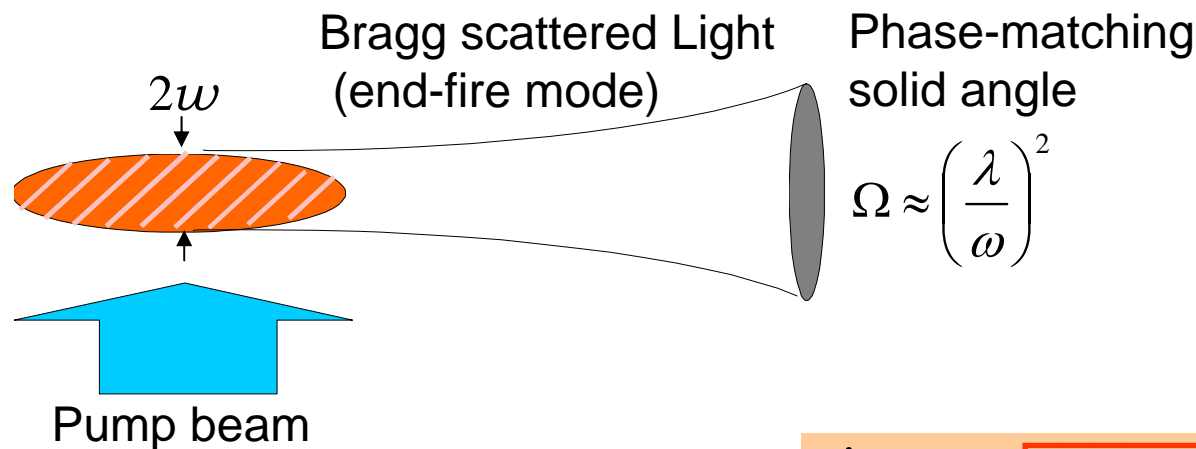
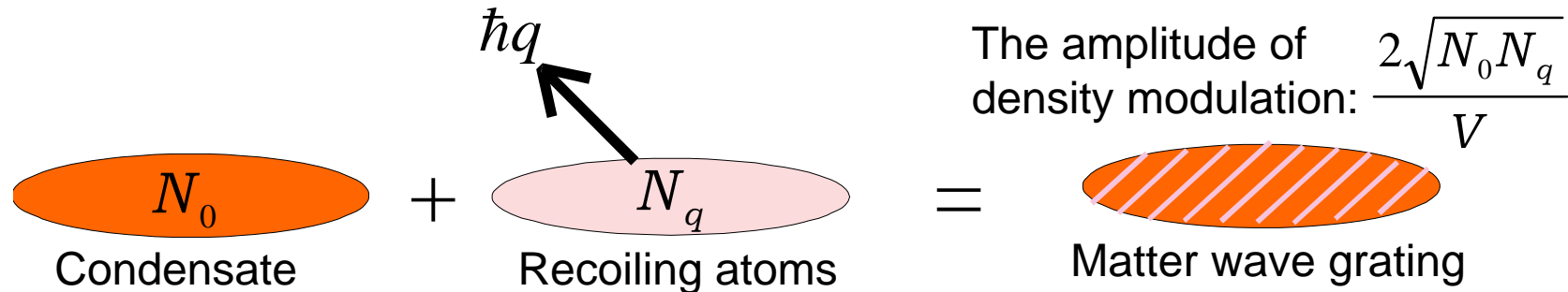
$$\left(g_0 \equiv \sqrt{\frac{d_{eg}^2 \omega}{2\epsilon_0 \hbar V}}, \quad d_{eg} \equiv \langle e | -e\hat{x} | g \rangle, \quad 2\kappa = \frac{1}{\tau_c} = \frac{\pi c}{lF}, \quad V = \frac{\pi}{4} \omega_0^2 \cdot l \right)$$

The ratio between R and spontaneous emission rate Γ (Purcell factor)

$$\frac{R}{\Gamma} = \frac{2g_0^2}{\kappa\Gamma} = 2C = \frac{3\lambda^3}{4\pi^2} \left(\frac{Q}{V} \right) \quad \left(Q \equiv \frac{\omega}{\Delta\omega} = \frac{2l}{\lambda} F \right)$$

Cooperativity parameter: $C \equiv \frac{g_0^2}{\kappa\Gamma} = \frac{12F}{\pi\omega_0^2 k^2} = \frac{F}{2\pi} \frac{\sigma_{\text{atom}}}{A} \quad \left(\sigma_{\text{atom}} = 6\pi\lambda^2, A = \frac{\pi}{4} \omega_0^2 \right)$

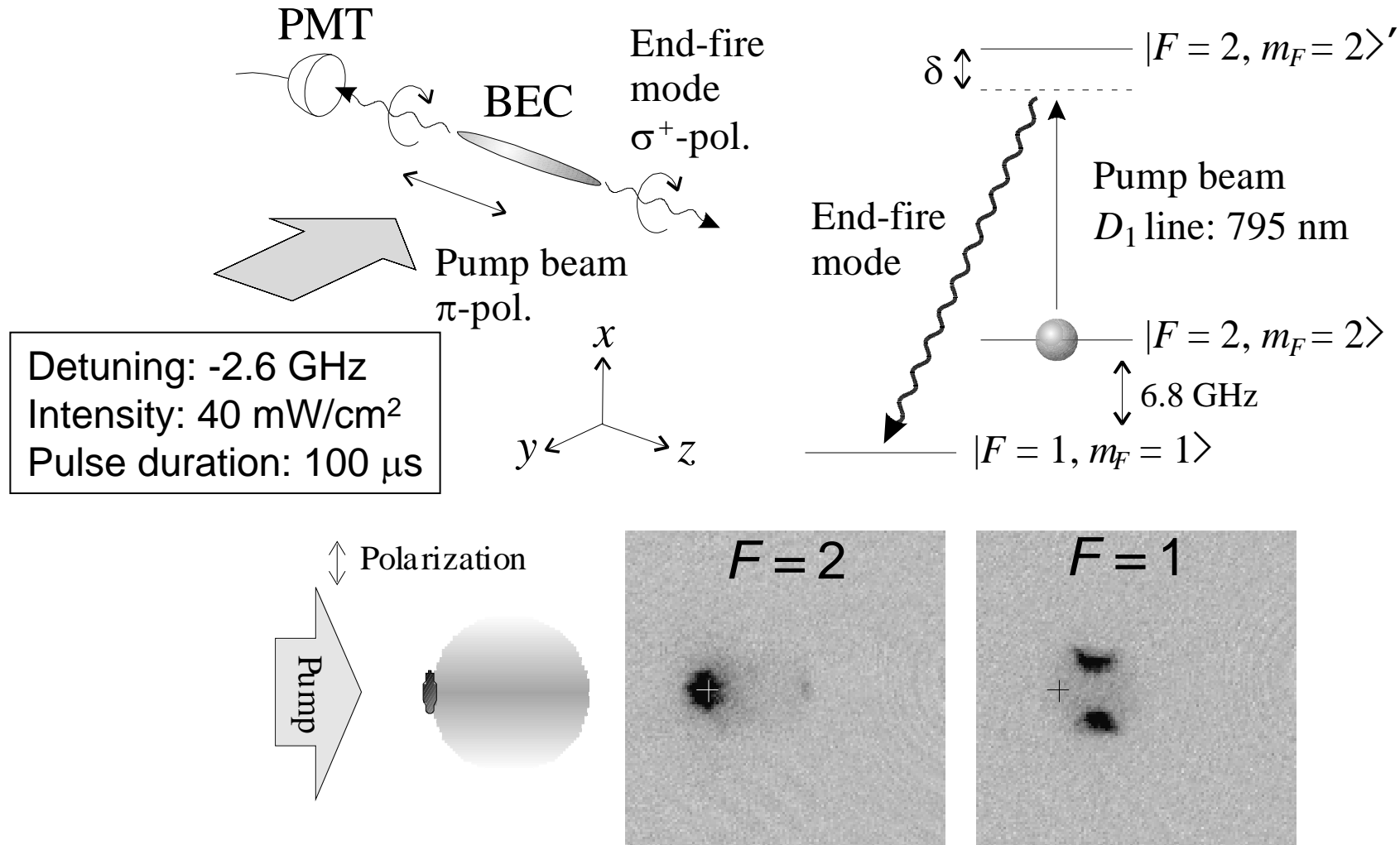
Relation between cooperativity parameter and superradiance



Power in the end-fire mode $P = \hbar\omega \frac{\sin^2 \theta}{8\pi/3} R N_0 N_q \Omega \rightarrow \frac{\dot{N}_q}{N_q} = R \frac{\sin^2 \theta}{8\pi/3} N_0 \Omega$

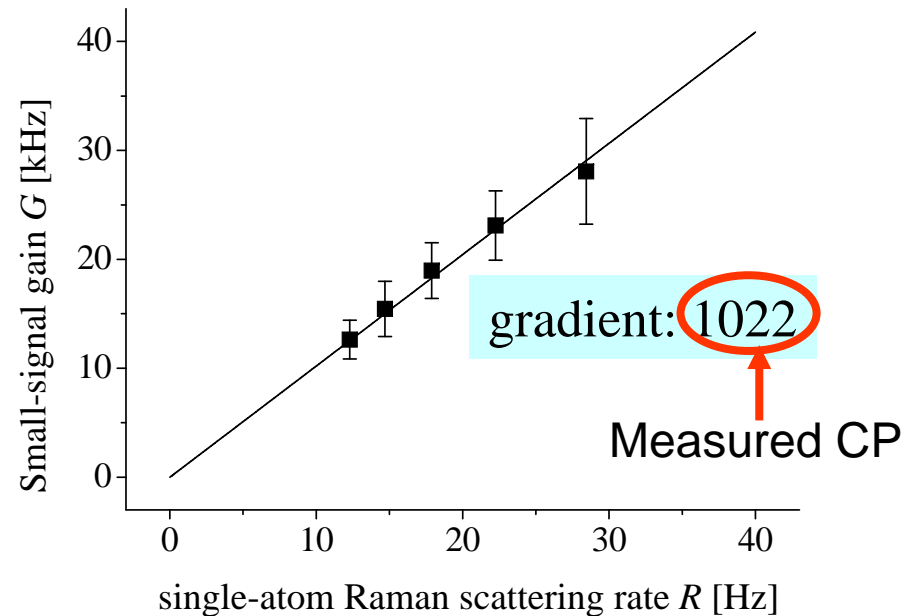
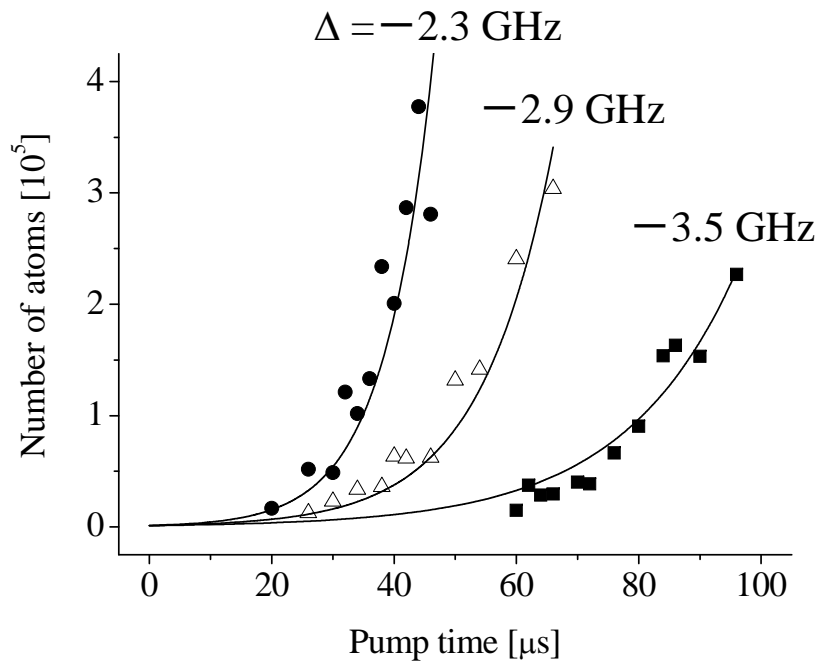
R : single-atom Rayleigh scattering rate Cooperativity parameter

Superradiant Raman scattering in a Bose condensate



Y. Yoshikawa, T. Sugiura, Y. T., and T. Kuga, PRA **69** 041603 (2004)

Measurement of the cooperativity parameter (CP) of a Bose condensate



$$\dot{N}_q \approx \frac{3}{8\pi} R N_0 N_q \Omega \xrightarrow{N_0 \ll N_q} N_q \approx e^{Gt}$$

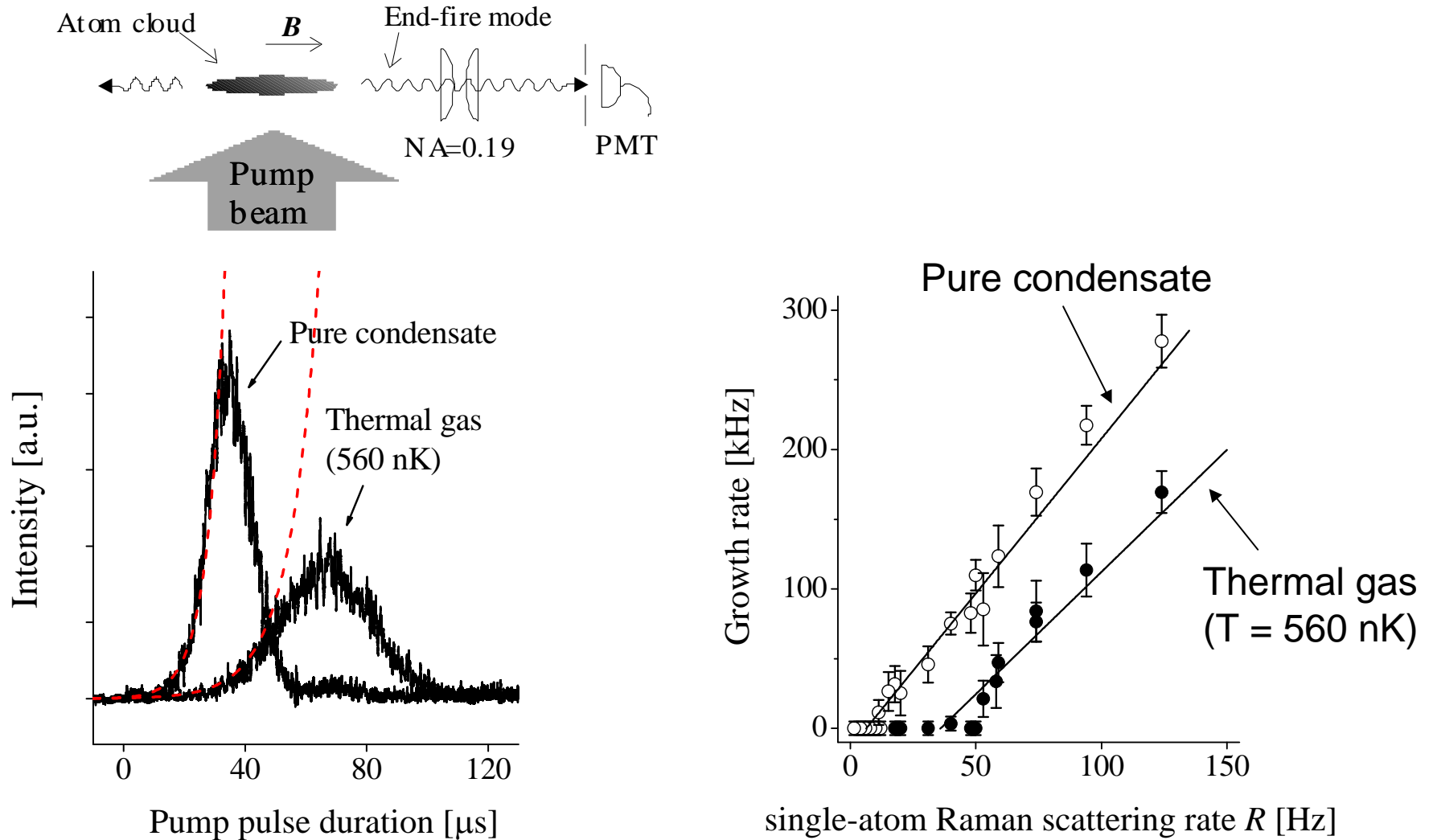
R : single-atom Raman scattering rate

Small-signal gain:

$$G = \frac{3}{8\pi} N_0 \Omega R = 890 R$$

Calculated CP

Superradiance in a Thermal gas



Y. Yoshikawa, Y. T. and T. Kuga, PRL **94** 083602 (2005)

What determines the coherence time?

a) The endfire mode

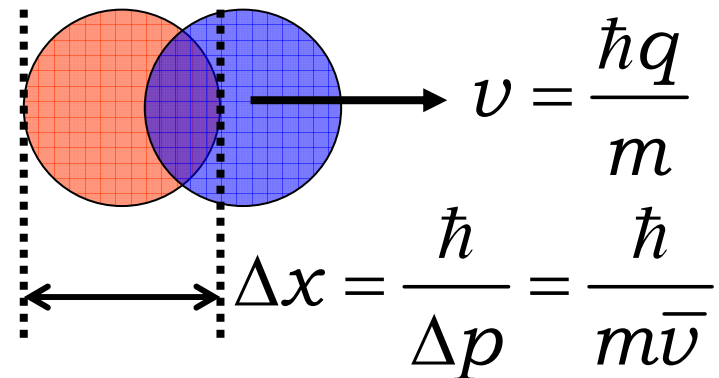
Doppler width:

$$\Delta\omega_D = q\bar{v} \quad \left(\bar{v} = \sqrt{\frac{k_B T}{m}} \right)$$

RMS velocity

$$\tau_c = \frac{1}{\Delta\omega_D} = \frac{1}{q\bar{v}}$$

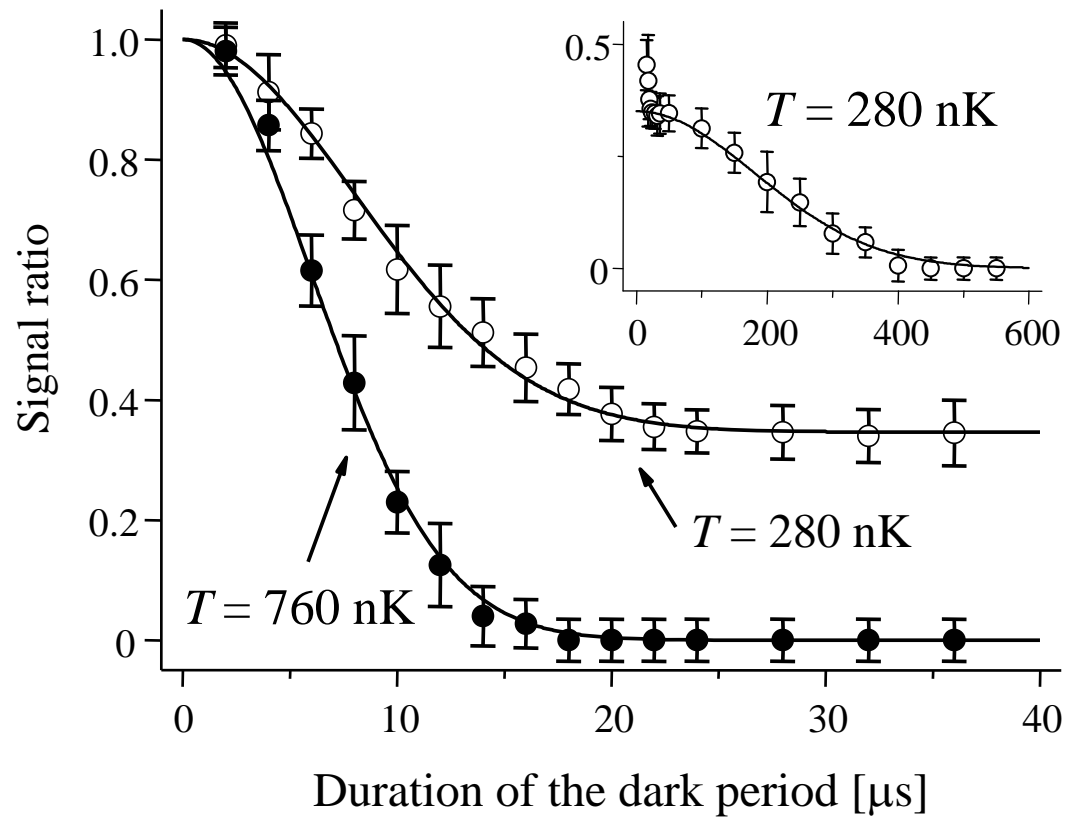
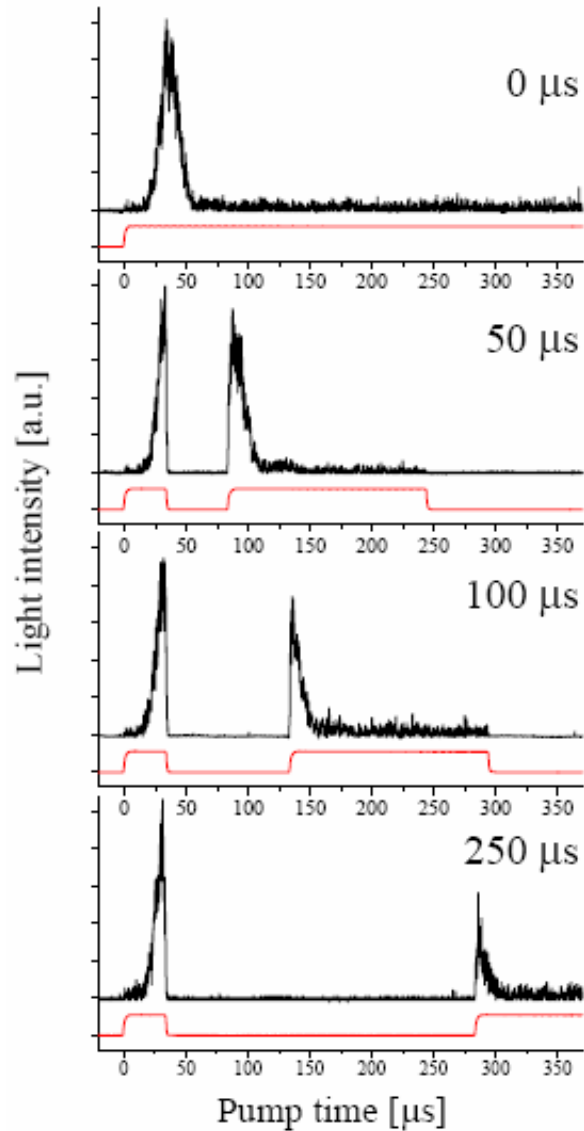
b) matter wave grating
(overlap of the wave packets)



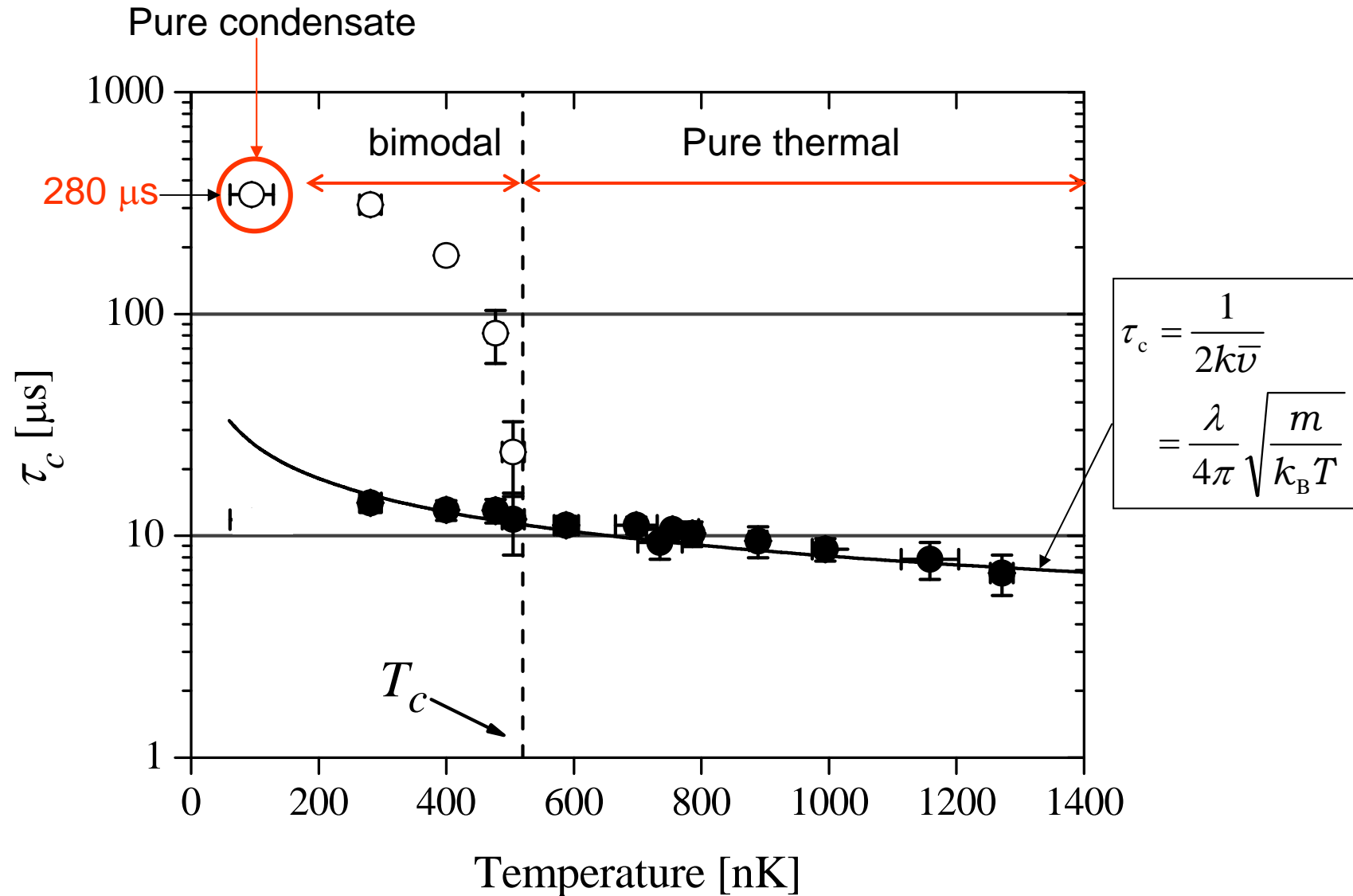
$$\tau_c = \frac{\Delta x}{v} = \frac{1}{q\bar{v}}$$

Coherence time is given by the inverse of the Doppler width

Storage (coherence) time measurement



Storage time vs. temperature



Merits of using a BEC for single photon storage

- Large cooperativity parameter ($\sim 10^3$)
(nearly 100% conversion efficiency)
- Long storage (coherence) time ($\sim 300\mu\text{s}$)
(could be extended by Lamb-Dicke effect)
- Arbitrary angle between the pump and the signal light (possibility of simultaneous storage of many photons)